XV. On the Thermal Effects of Fluids in Motion.—Part II.

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 ${
m In}$ the last experiment related in our former paper*, in which a low pressure of air was employed, a considerable variation of the cooling effect was observed, which it was necessary to account for in order to ascertain its influence on the results. therefore continued the experiments at low pressures, trying the various arrangements which might be supposed to exercise influence over the phenomena. already interposed a plug of cotton wool between the iron and copper pipes, which was found to have the very important effect of equalizing the pressure, besides stopping any solid or liquid particles driven from the pump, and which has therefore been retained in all the subsequent experiments. Another improvement was now effected by introducing a nozle constructed of boxwood, instead of the brass one previously used. This nozle is represented by fig. 1. Plate IV., in which a a is a brass casting which bolts upon the terminal flange of the copper piping, b b is a turned piece of boxwood screwing into the above, having two ledges for the reception of perforated brass plates, the upper plate being secured in its place by the turned boxwood c, which is screwed into the top of the first piece. The space enclosed by the perforated plates is 2.72 inches long and an inch and a half in diameter, and being filled with cotton, silk, or other material more or less compressed, presents as much resistance to the passage of the air as may be desired. A tin can d, filled with cotton wool, and screwing to the brass casting, serves to keep the water of the bath from coming in contact with the boxwood nozle.

In the following experiments, made in order to ascertain the variations in the cooling effect above referred to, the nozle was filled with 382 grs. of cotton wool, which was sufficient to keep up a pressure of about 34 lbs. on the inch in the tubes, when the pump was working at the ordinary rate. By opening the stopcock in the main pipe this pressure could be further reduced to about 22 lbs. by diminishing the quantity of air arriving at the nozle. By shutting and opening the stopcock we had therefore the means of producing a temporary variation of pressure, and of investigating its effect on the temperature of the air issuing from the nozle. In the first experiments the stopcock was kept open for a length of time, until the temperature of the rushing air became pretty constant; it was then shut for a period of $3\frac{3}{4}$, $7\frac{1}{2}$, 15, 30 or 60 seconds, then reopened. The oscillations of temperature thus pro-

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duced are laid down upon the Chart No. 1, in which the ordinates of the curves represent the temperatures according to the scale of thermometer C, each division corresponding to 0.0477 of a degree Centigrade. The divisions of the horizontal lines represent intervals of time equal to a quarter of a minute. The horizontal black lines show the temperature of the bath in each experiment.

The effect upon the pressure of the air produced by shutting the stopcock during various intervals of time, is given in the following Table:—

The last column gives also the effect occasioned by the permanent shutting or opening of the stopcock, 33.41 lbs. being nearly equal to the pressure when the stopcock has been closed for a long time.

In the next experiments, the opposite effect of opening the stopcock was tried, the results of which are laid down on Chart No. 2.

The effect upon the pressure of the air produced by opening the stopcock during the various intervals of time employed in the experiments, is exhibited in the next Table:—

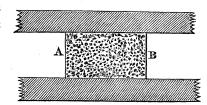
Stopcock opened for	$3\frac{3}{4}^{s}$.	7½s.	15s.	30s.	1 ^m .
Stopcock opened formsInitial pressurePressure after0 $3\frac{3}{4}$ Pressure after0 $7\frac{1}{2}$ Pressure after015Pressure after045Pressure after10Pressure after115Pressure after1130Pressure after130Pressure after145	34·37 29·57 32·47 33·5 33·94 34·1 34·2 34·33 34·37	$7\frac{1}{2}$ 5. $34\cdot 37$ $29\cdot 57$ $27\cdot 43$ $30\cdot 41$ $32\cdot 47$ $33\cdot 5$ $34\cdot 1$ $34\cdot 3$ $34\cdot 37$	34·37 29·57 27·43 25·15 30·41 32·4 33·5 33·94 34·14 34·30	34·37 29·57 27·43 25·15 23·23 29·4 32·13 33·24 33·90 34·14	34·37 29·57 27·43 25·15 23·23 22·76 28·82 31·44 32·9
Pressure after 2 0 Pressure after 2 15			34.37	34•33 34•37	33·66 34·06
Pressure after	•••••	•••••	•••••	34.37	34·06 34·20 34·37
Tressure andr 2 45	•••••		•••••	••••	0107

The remarkable fluctuations of temperature in the issuing stream accompanying such changes of pressure, and continuing to be very perceptible in the different cases

for periods of from 3 or 4 minutes up to nearly half an hour after the pressure had become sensibly uniform, depend on a complication of circumstances, which appear to consist of (1) the change of cooling effect due to the instantaneous change of pressure; (2) a heating or cooling effect produced instantaneously by compression or expansion in all the air flowing towards and entering the plug, and conveyed through the plug to the issuing stream; and (3) heat or cold communicated by contact from the air on the high-pressure side, to the metals and boxwood, and conducted through them to the issuing stream.

The first of these causes may be expected to influence the issuing stream instantaneously on any change in the stopcock; and after fluctuations from other sources have ceased, it must leave a permanent effect in those cases in which the stopcock is permanently changed. But after a certain interval the reverse agency of the second cause, much more considerable in amount, will begin to affect the issuing stream, will soon preponderate over the first, and (always on the supposition that this convection is uninfluenced by conduction of any of the materials) will affect it with all the variations, undiminished in amount, which the air entering the plug experiences, but behind time by a constant interval equal to the time occupied by as much air as is equal in thermal capacity to the cotton of the plug, in passing through the apparatus*; this, in the experiments with the stopcock shut, would be very exactly a

* To prove this, we have only to investigate the convection of heat through a prismatic solid of porous material, when a fluid entering it with a varying temperature is forced through it in a continuous and uniform stream. Let A B be the porous body, of length a and transverse section S; and let a fluid be pressed continuously through it in the direction from A to B, the temperature of this fluid as it enters at A being an arbitrary function F(t) of the time. Then if v be the



common temperature of the porous body and fluid passing through it, at a distance x from the end A, we have

$$\sigma \frac{dv}{dt} = k \frac{d^2v}{dx^2} - \frac{\theta}{S} \frac{dv}{dx}; \qquad (1)$$

if k be the conducting power of the porous solid for heat (the solid surrounding it being supposed to be an infinitely bad conductor, or the circumstances to be otherwise arranged, as is practicable in a variety of ways, so that there may be no lateral conduction of heat), σ the thermal capacity of unity of its bulk, and θ the thermal capacity of as much of the fluid as passes in the unit time. Now if, as is probably the case in the actual circumstances, conduction through the porous solid itself is insensible in its influence as compared with the convection of the fluid, this equation will become approximately

$$\sigma \frac{dv}{dt} = -\frac{\theta}{S} \frac{dv}{dx}, \qquad (2)$$

which, in fact, expresses rigorously the effect of the second cause mentioned in the text if alone operative. If F denote any arbitrary function, and if θ be supposed to be constant, the general integral of this equation is—

and if the arbitrary function be chosen to express by F(t) the given variation of temperature where the fluid enters the porous body, we have the particular solution of the proposed problem. We infer from it that, at any distance x in the porous body from the entrance, the temperature will follow the same law and extent of

quarter of a minute; but it appears to have averaged more nearly one-third of a minute in the varying circumstances of the actual experiments, since our observations (as may be partially judged from the preceding charts) showed us with very remarkable sharpness, in each case about twenty seconds after the shutting or opening of the stopcock, the commencement of the heating or cooling effect on the issuing stream, due to the sudden compression or rarefaction instantaneously produced in the air on the other side of the plug.

The entering air will, very soon after its pressure ceases to vary, be reduced to the temperature of the bath by the excellent conducting action of the spiral copper pipe through which it passes; and, consequently, twenty seconds or so later, the issuing stream can experience no further fluctuations in temperature except by the agency depending on the third cause.

That the third cause may produce very considerable effects is obvious, when we think how great the variations of temperature must be to which the surfaces of the solid materials in the neighbourhood of the plug on the high-pressure side are subjected during the sudden changes of pressure: and that the heat consequently taken in or emitted by these bodies may influence the issuing stream perceptibly for a quarter or a half hour after the changes of pressure from which it originated have ceased, is quite intelligible on account of the slowness of conduction of heat through the wood and metals, when we take into account the actual dimensions of the parts of the apparatus round the plug. It is not easy, however, to explain all the fluctuations of temperature which have been observed after the pressure had become constant in the different cases. Those shown in the first set of diagrams are just such as might be expected from the alternate heating and cooling which the solids must have experienced at their surfaces on the high-pressure side, and which must be conducted through so as to affect the issuing stream after a considerable time; but the great

variation as at the entrance, only later in time by an interval equal to $\frac{\sigma S}{\theta} x$. We conclude that the variations of temperature in the issuing stream due to the second cause alone, in the actual circumstances, are equal and similar to those of the air entering the plug, but later in time by $\frac{\sigma S a}{\theta}$. In this expression, the numerator, $\sigma S a$, denotes simply the thermal capacity of the whole plug. The plug, in the actual experiments, having consisted of 382 grains of cotton, of which the thermal capacity is about 191 times that of a grain of water, and (when the stopcock was closed) the air having been pumped through at the rate, per second, of 50 grains, of which the capacity is twelve times that of a grain of water, the value of $\frac{\sigma S a}{\theta}$ must have been $\frac{191}{12}$ seconds, or about a quarter of a minute. When the stopcock was open, an unknown quantity of air escaped through it, and therefore the value of $\frac{\sigma S a}{\theta}$ must have been somewhat greater. The variation which the value of θ must have experienced when the stopcock was opened or closed in the course of an experiment, or even merely in consequence of the change of pressure following the initial opening or closing of the stopcock, makes the circumstances not such as in any of the cases to correspond rigorously to the preceding solution; which, notwithstanding, represents the general nature of the convective effect nearly enough for the explanation in the text.

elevations of temperature shown in the second set of diagrams, which correspond to cases when the pressure was temporarily or permanently diminished, are not, so far as we see, explained by the causes we have mentioned, and the circumstances of these cases require further examination.

When we had thus examined the causes of the fluctuations of temperature in the issuing air, the precautions to prevent their injurious effect upon the accuracy of the determinations of the cooling effect in the passage of air through the porous plug became evident. These were simply to render the action of the pump as uniform as possible, and to commence the record of observations only after one hour and a half or two hours had elapsed from the starting of the pump. The system then adopted was to observe the thermometers in the bath and stream of air, and the pressure-gauge every two minutes or minute and a half; the means of which observations are recorded in the columns of the Tables. In some instances the air previous to passing into the pump was transmitted through a cylinder which had been filled with quicklime. But since by previous use its power of absorbing water had been considerably deteriorated, a portion of the air was always transmitted through a LIEBIG tube containing asbestos moistened with sulphuric acid or chloride of zinc. The influence of a small quantity of moisture in the air is trifling, but will hereafter be examined. That of the carbonic acid contained by the atmosphere was, as will appear in the sequel, quite inappreciable. It will be proper to observe that the thermometers by which the temperature of the bath and issuing air was ascertained, were repeatedly compared together to avoid any error which might arise from the alteration of their fixed points from time to time.

Table I	-Ex	periments	with a	plug	consisting	of	191	grains of	cotton	wool.
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1.	2.	3.	4.	5.	6.	7.	8.
Number of observations from which the results in Columns 4, 6, and 7, are obtained.	Cubic inches	Water in 100 grains of air, in grains.	Pressure in lbs. on the square inch.	Atmospheric pressure.	Temperature of the bath.	Temperature of the issuing air.	Cooling effect in Cent. degrees.
20 20 10 10 10 10	10822 10998 Not observed. 10769 10769 10769	0·51 0·30 0·56 0·66 0·66 0·66	21·326 21·239 20·446 20·910 20·934 20·995 20·933	14·400 14·252 14·609 14·772 14·775 14·779	20·295 16·740 17·738 16·039 16·065 16·084 16·081	20·201 16·615 17·622 15·924 15·967 15·984 15·974	0.094 0.125 0.116 0.115 0.098 0.100 0.107
Mean	••••	0.57	20.969	14.624	17.006	16.898	0.108

In each, excepting the first of the seven experiments above recorded, the air was passed through the quicklime cylinder.

In the next experiments the nozle was filled with 382 grains of cotton wool. The

intermediate stopcock was however partly opened, in order that by discharging a portion of the air before its arrival at the nozle, the pressure might not be widely different from that employed in the last series. In all excepting the last experiment recorded in the following Table, the cylinder of lime was dispensed with.

Table II.—Experiments with a smaller quantity of air passed through a plug consisting of 382 grs. of cotton wool.

1.	2.	3.	4.	5.	6.	7.	8.
Number of observations from which the results in Columns 4, 6, and 7, are obtained.	Cubic inches passed through the nozle per minute.	Water in 100 grains of air, in grains.	Pressure in lbs. on the square inch.	Atmospheric pressure.	Temperature of the bath.	Temperature of the issuing air.	Cooling effect in Cent. degrees.
20 30	3865 3960	0·59 0·73	22·614 22·818	14·513 14·514	20°·363 19·853	20°224 19°769	°0·139 0·084
20	Not observed.	0.56	22.818	14.604	20.481	20.407	0.074
45	3125	0.65	22.296	14.590	20.584	20.313	0.271
20	Not observed.	1.23	23.000	14.518	18.636	18.476	0.160
36	Not observed.	1.20	22.616	14.520	20.474	20.336	0.138
50	Not observed.	1.36	22.582	14.518	20.485	20.325	0.160
Mean	••••	0.90	22.678	14.540	20.125	19.979	0.146

Table III.—Experiments in which the entire quantity of air propelled by the pump was passed through a plug consisting of 382 grains of cotton wool. The cylinder of lime was not employed.

1.	2.	3.	4.	5.	6.	7.	8.
Number of observations from which the results in Columns 4, 6, and 7, are obtained.	Cubic inches passed through the nozle per minute.	Water in 100 grains of air, in grains.	Pressure in lbs. on the square inch.	Atmospheric pressure.	Temperature of the bath.	Temperature of the issuing air.	Cooling effect in Cent. degrees.
7 10 10 10 10	11766 Not observed. Not observed. Not observed. Not observed.	0.56 0.56 0.36 0.36 0.36	36·625 35·671 35·772 35·872 36·026	14·583 14·790 14·504 14·504 14·504	19.869 20.419 16.096 16.104 16.232	19·535 20·098 15·730 15·721 15·869	0.334 0.321 0.366 0.383 0.363
Mean		0.44	35.993	14.577	17.744	17:390	0.354

In the next series of experiments the air was passed through a plug of silk, formed by rolling a silk handkerchief into a cylindrical shape, and then screwing it into the nozle. The silk weighed 580 grains, and the small quantity of cotton wool placed on the side next the thermometer in order to equalize the stream of air more completely, weighed 15 grains. The stopcock was partly opened as in the experiments of Table II., in order to reduce the pressure to that obtained by passing the full

quantity of air propelled by the pump through a more porous plug. The cylinder of lime was employed.

Table IV.—Experiments in which a smaller quantity of air was passed through a plug consisting of 580 grains of silk.

1.	2.	3.	4.	5.	6.	7.	8.
Number of observations from which the results in Columns 4, 6, and 7, are obtained.	Cubic inches	Water in 100 grains of air, in grains.	Pressure in lbs. on the square inch.	Atmospheric pressure.	Temperature of the bath.	Temperature of the issuing air.	Cooling effect in Cent. degrees.
10 10 10 10	3071 Not observed. Not observed. Not observed.	0·18 0·18 0·14 0·14	33·168 33·024 33·820 33·226	14·727 14·732 14·660 14·650	18.882 18.884 19.066 19.068	18·524 18·536 18·686 18·695	0.358 0.348 0.380 0.373
Mean	****	0.16	33.309	14.692	18.975	18.610	0.365

Table V.—Experiments in which the entire quantity of air propelled by the pump was passed through the silk plug. The cylinder of lime was employed in all excepting the first two experiments.

1.	2.	3.	4.	5.	6.	7.	8.
Number of observations from which the results in Columns 4, 6, and 7, are obtained.	Cubic inches passed through the nozle per minute.	Water in 100 grains of air, in grains.	Pressure in lbs. on the square inch.	Atmospheric pressure.	Temperature of the bath.	Temperature of the issuing air.	Cooling effect in Cent. degrees.
10	br to a	0.40	* 0 * 00	14.500	2 2 2 2	2000	9 0
10	7594	0.40	53.722	14.580	17.585	1 <u>6</u> ·903	0.682
10	Not observed.	0.40	53.530	14.580	17.628	16.954	0.674
10	Not observed.	0.32	53.317	14.563	17.993	17.318	0.675
10	Not observed.	0.32	53.317	14.563	18.027	17.357	0.670
10	7742	0.11	55.797	14.615	17.822	17.063	0.759
10	Not observed.	0.11	54.074	14.611	17.813	17.079	0.734
10	Not observed	0.11	55.720	14.608	17.808	17.082	0.726
10	Not observed.	0.11	56.174	14.605	17.796	17.058	0.738
Mean		0.23	54.456	14.591	17.809	17.102	0.707

In order to obtain a greater pressure, a plug was formed of silk "waste" compressed very tightly into the nozle.

Table VI.—Experiments in which the a	air, after passing through the cylinder
of lime, was forced through a plug	consisting of 740 grains of silk.

1.	2.	3.	4.	5.	6.	7.	8.
Number of observations from which the results in Columns 4, 6, and 7, are obtained.	Cubic inches	Water in 100 grains of air, in grains.	Pressure in lbs. on the square inch.	Atmospheric pressure.	Temperature of the bath.	Temperature of the issuing air.	Cooling effect in Cent. de- grees.
1.0	NT . 1	0.10	* 0.0*0		2	0,	0 00
10	Not observed.		79.852	14.777	17.050	15.884	1.166
10	Not observed.		80.133	14.782	17.066	15.913	1.153
10	Not observed.		79.870	14.787	17.079	15.945	1.134
10	5650	0.19	80.013	14.793	17.083	15.967	1.116
10	Not observed.	0.15	79.814	14.960	16.481	15.338	1.143
10	Not observed.	0.15	80.274	14.957	16.489	15.374	1.115
10	Not observed.	0.15	79.903	14.953	16.505	15.392	1.113
10	5378	0.15	77.867	14.950	16.521	15.428	1.093
10	Not observed.	0.14	78.214	14.638	12.851	11.770	1.081
10	Not observed.	0.14	78.245	14.638	12.877	11.800	1.077
10	Not observed.	0.14	78.180	14.638	12.885	11.824	1.061
10	Not observed.	0.14	78.633	14.638	12.905	11.839	1.066
Mean		0.16	79.250	14.793	15.483	14.373	1.110

In the foregoing experiments the pressure of the air on its exit from the plug was always exactly equal to the atmospheric pressure. To ascertain the effect of an alteration in the pressure of the exit air, we now enclosed a long siphon barometer within the glass tube (fig. 10). The upper part of this tube was surmounted with a cap, furnished with a stopcock, by partially closing which the air at its exit could be brought to the required pressure. The influence of pressure in raising the mercury in the thermometer by compressing its bulb, was ascertained by plunging the instrument into a bottle of water within the glass tube, and noting the amount of the sudden rise or fall of the quicksilver on a sudden augmentation or reduction of pressure. It was found that the pressure equal to that of 17 inches of mercury, raised the indication by $0^{\circ}.09$; which quantity was therefore subtracted after the usual reduction of the thermometric scale.

Table VII.—Experiments with the plug consisting of 740 grains of silk. Pressure of the exit air increased. Cylinder of lime used.

1.	2.	3.	4,	5.	6.	7.	8.
Number of observations from which the results in Columns 4, 6, and 7, are obtained.	Cubic inches passed through the nozle per minute.	Water in 160 grains of air, in grains.	Pressure in lbs. on the square inch.	Pressure of the exit air.	Temperature of the bath.	Temperature of the issuing air.	Cooling effect in Cent. degrees.
10 10 10 10	Not observed. Not observed. Not observed. Not observed.	0.14	82·982 82·510 81·895 80:630	23·093 22·878 22·798 22·488	12.673 12.713 12.755 12.795	11.612 11.676 11.725 11.792	1.061 1.037 1.030 1.003
Mean	Estimated at 5400	0.14	82.004	22.814	12:734	11.701	1.033

With reference to the experiments in Table VII. it may be remarked, that the cooling effect must be the excess of that which would have been obtained had the air been only resisted by the atmospheric pressure in escaping from the plug, above the cooling effect that would be found in an experiment with the temperature of the bath and the pressure of the entering air the same as the temperature and pressure of the exit air in the actual experiment, and the air issuing at atmospheric pressure. Hence, since two or three degrees of difference of temperature in the bath would not sensibly alter the cooling effect in any of the experiments on air, the cooling effect in an experiment in which the pressure of the exit air is increased, must be sensibly equal to the difference of the cooling effects in two of the ordinary experiments, with the high pressures the same as those used for the entering and issuing air respectively, and the low pressure that of the atmosphere in each case; a conclusion which is verified by the actual results, as the comparison given below shows.

The results recorded in the foregoing Tables are laid down on Chart No. 3, in which the horizontal lines represent the excess of the pressure of the air in the receiver over that of the exit air as found by subtracting the fifth from the fourth columns of the Tables, and the vertical lines represent the cooling effect in tenths of a degree Centigrade. It will be remarked that the line drawn through the points of observation is nearly straight, indicating that the cooling effect is, approximately at least, proportional to the excess of pressure, being about 018° per pound on the square inch of difference of pressure. Or we may arrive at the same conclusion by dividing the cooling effect (δ) by the difference of pressures (P-P') in the different experiments. We thus find, from the means shown in the different tables,—

=:0170
.0179
.0165
.0196
.0177
.0172
.0174
.0176

On the Cooling Effects experienced by Carbonic Acid in passing through a porous Plug.

The position of the apparatus gave us considerable practical facilities in experimenting with carbonic acid. A fermenting tun 10 feet deep and 8 feet square was filled with wort to a depth of 6 feet. After the fermentation had been carried on for about forty hours, the gas was found to be produced in sufficient quantity to supply the pump for the requisite time. The carbonic acid was conveyed by a gutta-percha pipe, and passed through two glass vessels surrounded by ice in order to condense the greater portion of vapours. In the succeeding experiment the total quantity of liquid so condensed was 300 grains, which having a specific gravity of 9965, was composed of 10 grains of alcohol and 290 grains of water. On analysing a portion of the gas during the experiment by passing it through a tube containing chloride of zinc, it was found to contain 0.733 gr. of water to 100 grs. of carbonic acid.

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Table VIII.—Carbonic acid forced through a plug of 382 grs. of cotton wool. Mean barometric pressure 29.45 inches, equivalent to 14.399 lbs. Gauge under atmospheric pressure 151. The pump was placed in connexion with the pipe immersed in carbonic acid at 10^h 55^m.

1.	2.	3,	4.	5.	6.
Time of observa-	Volume percentage of carbonic acid.	Pressure-gauge; mean pressure in lbs. on the square inch.	Temperature of the bath, by indications of thermometer. Temperature of the issuing gas, by indications of thermometer.		Cooling effect in Cent. degrees.
h m 10 47	0	79.0	486•0	198•5	
49	0	79.0	486.0	198.5	
53	0	79.6	486.0	198.2	
57		85.2	486.0	195.0	
58		86.0	486.0	186.0	
5 9)	85.07	486.07	188∙6 ີ	
11 0	95.51	85.0	486.0	188.5	
2		86.4	486.0	187.6	
4	İ	86.7	486.0	187.8	
6	1	86.6	486.0	188.9	
9	95.51	86.6	486.0	188.9	
13		84.0	486.0	188.65	
14	07 77 04 00	84.2 lbs.	486.0	188.1	0
15	95.51 >94.89	84.4 > 84.906 = 32.989	$ 486.0\rangle 486.00 = 20.001$	188.0 > 188.36 = 18.611	1.390
19		84.5	486·0 486·0	188.0	
22 24		84·1 84·6	486.0	188·1 188·3	ĺ
$\frac{24}{25}$	93.03	84.2	486.0	188.5	}
28 28	90 00	84.1	486.0	188.6	
32	·	83.2	486.0	188.9	
33	1 3	83.87	486.0	188.9	
35	86.82	84.0	486.0	189.0	
40		83.8	486.0	189.6	
41		83.9	486.0	189.7	į ,
43		85.0	485.9	189.9	
45	79.37 >80.61	86.0 > 84.245 = 33.286	485.9 >485.94=19.998	190.4 > 190.1 = 18.787	1.211
49		84.6	485.9	190.8	
51		84.5	485.9	190.8	
53		83.9	485.9	190.6	
55	75.65	83.6	485.9	190.6	
12 0		83.6	485.9	190.8	
2	70.68	83.0	485.7	190.8	
5	1 1	82.7	485.7	190.9	
$\begin{array}{c} 9 \\ 13 \end{array}$	68.82	$\begin{vmatrix} 82.7 \\ 82.9 \end{vmatrix} $	$\begin{vmatrix} 485.4 \\ 485.4 \end{vmatrix}$ $\begin{vmatrix} 485.52 = 19.980 \end{vmatrix}$	$\begin{vmatrix} 190.8 \\ 191.1 \end{vmatrix}$ \Rightarrow 191.07 = 18.884	1.096
15	66.96	82.7	485.5	191.3	
21	30 00	82.7	485.4	191.5	
23	\	82.85	485.47	191.55	
25	65.72	82.9	485.4	191.6	
28		82.9	485.4	191.7	
33		82.2	485.4	191.8	
35	63.23	82.3	485.4	191.7	
40		81.9	485.3	191.65	
44	00.00	81.9	485.2	191.6	
45	63.23 >63.85		$ 485\cdot2 $ $>485\cdot18=19\cdot966$		1.007
52	60.0	82.4	485.0	191.65	
55	62.0	83.9	485.0	192.0	
$\begin{array}{cc} 1 & 2 \\ & 5 \end{array}$	63.23	84.1	485.0	192.1	
11	03.29	84.9	485·0 485·0	192.1	
15	65.72	82.1	484.9	192·3 192·1	
10	10012	ر ۱	1019)	13% 1	

Table IX.—Carbonic acid forced through a plug consisting of 191 grs. of cotton wool. Mean barometric pressure 29.6 inches, equivalent to 14.472 lbs. Gauge under atmospheric pressure 150.6. Pump placed in connexion with the pipe immersed in carbonic acid at 10^h 38^m.

1.	2.	3.	4.	5.	6.
Time of observation.	Volume per- centage of car- bonic acid.	Pressure-gauge, and pressure in lbs. on the square inch equivalent thereto.	Indication of thermometer. Temperature of the bath.	Indication of thermometer. Temperature of the issuing gas.	Cooling effect in Cent. degrees.
h m 10 40 42 44 50 53 55 57 59 11 0 1 3 5 7 9 10 11 15 17 19 20 21 25	95.51	122.9 122.6 122.6 122.6 122.5 122.8 122.1	$ \begin{array}{c} 461 \cdot 5 \\ 461 \cdot 6 \\ 461 \cdot 75 \\ 461 \cdot 75 \\ 461 \cdot 75 \\ 461 \cdot 75 \\ 461 \cdot 9 \\ 461 \cdot 9 \\ 461 \cdot 9 \\ 462 \cdot 0 \\ 462 \cdot 0 \\ 462 \cdot 0 \\ 462 \cdot 0 \\ 462 \cdot 2 \end{array} $	187·55 187·55 187·6 187·55 188·1 188·1	

In the above, as well as in the next series, the carbonic acid contained 0.35 per cent. of water.

Table X.—Experiment in which carbonic acid was forced through a plug consisting of 580 grs. of silk. Mean barometric pressure 29.56, equivalent to 14.452 lbs. Gauge under atmospheric pressure 150.8. Pump placed in connexion with the pipe immersed in carbonic acid at 12^h 53^m. Quantity of gas forced through the plug about 7170 cubic inches per minute.

1.	2.	3.	4.	5.	6.
Time of observation.	Volume per- centage of car- bonic acid.	Pressure-gauge, and pressure in lbs. on the square inch equivalent thereto.	Indication of thermometer. Temperature of the bath.	Indication of thermometer. Temperature of the issuing gas.	Cooling effect in Cent. degrees.
h s 12 42 44 46 49 50 52 54 57	0 0 0 0 0 0	52.2 52.2 52.2 52.2 52.2 52.2 52.2 56.0 55.7	$ \begin{vmatrix} 464 \cdot 2 \\ 464 \cdot 35 \\ 464 \cdot 4 \\ 464 \cdot 35 \\ 464 \cdot 35 \\ 464 \cdot 4 \\ 464 \cdot 55 \\ 464 \cdot 65 \end{vmatrix} $	$ \begin{vmatrix} 185.6 \\ 185.5 \\ 185.5 \\ 185.55 \\ 185.55 \\ 185.5 \\ 179.0 \\ 166.3 \end{vmatrix} $ $ 185.53 = 18.323$	ỏ ∙749
1 0 5 7 9 10 11	95.51	56·0 56·0 56·0 56·0 55·92=51·7	464·3 464·55 464·5 464·6 464·55 464·47=19·077	165.0 165.0 164.9 164.9 164.8	2.821
17 20 24 25 27 30 35 36	93·03 85·92	56·0 55·5 55·7 56·1 56·0 55·94=51·68 56·1 56·1 56·1	464.5 464.6 464.6 464.6 464.7 464.7 464.7 464.8 464.8 464.9	165.0 165.4 166.0 166.3 166.8 167.9 168.9 169.1 169.6	2.550

In the above experiment, as well as in those of the adjoining Tables, the sudden diminution of pressure on connecting the pump with the receiver containing carbonic acid, is in perfect accordance with the discovery by Professor Graham of the superior facility with which that gas may be transmitted through a porous body compared with an equal volume of atmospheric air.

Table XI.—Experiment in which carbonic acid was forced through a plug consisting of 740 grs. of sik. Mean barometric pressure 30.065, equivalent to 14.723 lbs. on the inch. Gauge under atmospheric pressure 145.65. Pump placed in connexion with the pipe immersed in carbonic acid at 11^h 37^m. Per-centage of moisture in the carbonic acid 0.15.

1.	2.	3.	4.	5.	6.
Time of observation.	Volume per- centage of car- bonic acid.	Pressure-gauge, and pressure in lbs. on the square inch equivalent thereto.	Indication of thermometer. Temperature of the bath.	Indication of thermometer. Temperature of the issuing gas.	Cooling effect in Cent. degrees.
h m 11 28 30 32 34 36 37 38 39 43 45 47 50 53 55 57 12 0 2 5	95·51 95·51 95·51 95·51 95·51 95·51	35·5 35·1 35·6 35·2 35·2 35·2 36·0 36·2 36·6 36·9 37·0 37·0 37·0 37·0 37·0 37·0 37·0 37·0	318·9 318·95 318·95 318·95 318·95 318·95 318·95 319·03 319·05 319·15 319·3 319·3	117.9 118.0 118.0 118.0 117.9 117.73 117.5 112.0 94.0 93.95 83.6 83.0 82.6 82.4 82.35 82.3 82.7 83.0 83.0 83.0	å·87

In order to ascertain the cooling effect due to pure carbonic acid, we may at present neglect the effect due to the small quantity of watery vapour contained by the gas; and as the cooling effects observed in the various mixtures of atmospheric air and carbonic acid appear nearly consistent with the hypothesis that the specific heats of the two elastic fluids are for equal volumes equal to one another, and that each fluid experiences in the mixture the same absolute thermo-dynamic effect as if the other were removed, we may for the present take the following estimate of the cooling effects due to pure carbonic acid, at the various temperatures and pressures employed, calculated by means of this hypothesis from the observations in which the per-centage of carbonic acid was the greatest, and in fact so great, that a considerable error in the correction for the common air would scarcely affect the result to any sensible extent.

	Temperature of the bath.	Excess of pressure, P-P'.	Cooling effect, δ .	Cooling effect divided by excess of pressure.
From Table IX From Table VIII From Table X From Table XI	20·001 19·077	5•958 18•590 37•248 60•601	0·459 1·446 2·938 5·049	.0770 .0778 .0789 .0833
	Mean 17·721			Mean of first three .0779
				Mean of all0793

We shall see immediately that the temperature of the bath makes a very considerable alteration in the cooling effect, and we therefore select the first three results, obtained at nearly the same temperature, in order to indicate the effect of pressure. On referring to Chart No. 3, it will be remarked that these three results range themselves almost accurately in a straight line. Or, by looking to the numbers in the last column, we arrive at the same conclusion.

Cooling Effect experienced by Hydrogen in passing through a porous Plug.

Not having been able as yet to arrange the large apparatus so as to avoid danger in using this gas in it, we have contented ourselves for the present with obtaining a determination by the help of the smaller force-pump employed in our preliminary experiments. The hydrogen, after passing through a tube filled with fragments of caustic potash, was forced, at a pressure of 68.4 lbs. on the inch, through a piece of leather in contact with the bulb of a small thermometer, the latter being protected from the water of the bath by a piece of india-rubber tube. At a temperature of about 10° Cent., a slight cooling effect was observed, which was found by repeated trials to be 0°.076. The pressure of the atmosphere being 14.7 lbs., it would appear that the cooling effect experienced by this gas is only one-thirteenth of that observed with atmospheric air. We state this result with some reserve, on account of the imperfection of such experiments on a small scale, but there can be no doubt that the effect of hydrogen is vastly inferior to that of atmospheric air.

Influence of Temperature on the Cooling Effect.

By passing steam through pipes plunged into the water of the bath, we were able to maintain it at a high temperature without any considerable variation. The passage of hot air speedily raised the temperature of the stem of the thermometer, as well as of the glass tube in which it was enclosed; but nevertheless the precaution was taken of enclosing the whole in a tin vessel, by means of which water in constant circulation with the water of the bath was kept within one or two inches of the level of the mercury in the thermometer. The bath was completely covered with a wooden lid, and the water kept in constant and vigorous agitation by a proper stirrer.

Table XII.—Experiment in which—1st, air; 2nd, carbonic acid; 3rd, air dried by quicklime was forced through a plug consisting of 740 grs. of silk. Mean barometric pressure 30.015, equivalent to 14.68 lbs. on the inch. Gauge under the atmospheric pressure 150. Per-centage of moisture in the carbonic acid 0.31. Pump placed in connexion with the pipe immersed in carbonic acid at 11^h 24^m. Disconnected and attached to the quicklime cylinder at 12^h 22^m.

1.	2.	3.	4.	5.	6.
Time of observation.	Volume per- centage of car- bonic acid.	Pressure-gauge, and pressure in lbs. on the square inch equivalent thereto.	Indication of thermometer. Temperature of the bath.		
h m 11 5 7 9 11 13	0 0 0 0 0	31.6 31.4 31.7 31.6 31.9 31.5 31.62=91.508	646·35 646·3 646·1 646·05 646·05 646·05	479·1 478·8 478·05 478·1 478·2 478·35 478·43=90·008	°1•444
17 19 21 22 23 24 25	0 0 0 0 0 0	31.8 31.5 32.0 32.2 32.2 32.0 32.0	$ \begin{array}{c c} 646 \cdot 2 \\ 646 \cdot 0 \\ 646 \cdot 0 \\ 646 \cdot 1 \\ 646 \cdot 1 \\ 646 \cdot 1 \\ 646 \cdot 1 \end{array} $ $ \begin{array}{c c} 646 \cdot 08 = 91 \cdot 442 \\ 646 \cdot 08 = 91 \cdot 442$	478.7 478.6 478.7 478.6 478.1 478.8 477.0	1·399
26 30 32 33 36 38 40 43	$ \begin{vmatrix} 0 \\ 95.51 \\ 95.51 \\ 95.51 \\ 95.51 \\ 95.51 \\ 93.03 \end{vmatrix} $	32·1 32·2 32·2 32·0 32·6 32·2 32·2 32·2 32·2	646.4 646.7 646.5 646.45 646.6 646.6 646.6	471.6 469.2 469.5 469.6 469.6 469.9 469.93 470.05	3·472
46 48 50 53 55 58	90.60 \ 91.81	$ \begin{vmatrix} 32 \cdot 1 \\ 32 \cdot 1 \\ 32 \cdot 1 \end{vmatrix} $ $ \begin{vmatrix} 32 \cdot 1 \\ 32 \cdot 1 \end{vmatrix} $ $ \begin{vmatrix} 32 \cdot 1 \\ 32 \cdot 0 \end{vmatrix} $ $ \begin{vmatrix} 32 \cdot 1 \\ 32 \cdot 0 \end{vmatrix} $ $ \begin{vmatrix} 22 \cdot 16 \\ 32 \cdot 0 \end{vmatrix} $	$ \begin{vmatrix} 647.0 \\ 647.1 \\ 647.4 \\ 647.2 \\ 647.2 \\ 647.2 \end{vmatrix} $ $ \begin{vmatrix} 647.03 = 91.579 \\ 647.2 \\ 647.2 \end{vmatrix} $	$ \begin{vmatrix} 470.3 \\ 470.9 \\ 471.05 \end{vmatrix} 470.57 = 88.255 $ $ 471.75 \\ 472.05 \end{vmatrix} 472.00 = 88.638$	3.324
12 0 4 6 9	75·65 77·37 75·65 65·72 65·72	32·0 32·6 32·25 32·8 32·4	647.7 647.9 647.8 647.95 647.95	472.6 472.6 472.9 473.25 473.95 474.1	3.009
15 20 22 27 29 31 33 35	60.83 62.46 60.83 62.46 60.83 62.46 60.83 62.46 60.83 62.46 60.83 62.46 60.83 62.46 60.83 62.46 60.83 62.46 60.83 62.46 60.83 62.46 60.83 62.46 60.83 62.46 60.83 62.46 60.83 62.46 60.83 62.46 60.83 62.46 60.83 60.83 62.46 60.83	32·2 32·4 32·9 32·0 31·6 32·0 32·1 32·2	647.95 647.95 647.85 647.8 647.8 647.3 647.1 647.0	474·8 475·15 475·2 477·0 480·1 480·6 480·6 480·8 480·8	2.549
39 41 43 45 47 49 51	0 0 0 0 0 0 0	32.0 32.1 32.1 32.1 32.1 32.1 32.2 32.3 =89.618 32.6 32.8	647·1 7 647·03 647·1	480·9 481·03 480·9 481·02 481·04 480·98 480·9	1.050

Although hot air had been passed through the plug for half an hour before the readings in the preceding Table were obtained, it is probable that the numbers 1.444 and 1.399, representing the cooling effect of atmospheric air, are not so accurate as the value 1.050. Taking this latter figure for the effect of an excess of pressure of 89.618—14.68=74.938 lbs., we find a considerable decrease of cooling effect owing to elevation of temperature, for that pressure, at the low temperatures previously employed, is able to produce a cooling effect of 1.309.

In order to obtain the effect of carbonic acid unmixed with atmospheric air, we shall, in accordance with the principle already adhered to, consider the thermal capacities of the gases to be equal for equal volumes. Then the cooling effect of the

pure gas=
$$\frac{3.472 \times 100 - 1.052 \times 4.49}{95.51}$$
=3°.586.

Collecting these results, we have,-

Temperature of bath.	Excess of pressure.	Cooling effect.	Cooling effect reduced to 100 lbs. pressure.	Theoretical cooling effect for 100 lbs. pressure.
12.844	60·601	Š·049	8 ·33	8∙27
19.077	37·248	2·938	7 · 89	8•07
91.516	74·938	3·586	4·78	4•96

Note.—The numbers shown in the last column of the Table are calculated by the general expression given in our former paper* for the cooling effect, from an empirical formula for the pressure of carbonic acid, recently communicated by Mr. Rankine in a letter, from which the following is extracted.

"Glasgow, May 9, 1854.

"Annexed I send you formulæ for carbonic acid, in which the coefficient a has been determined solely from Regnault's experiments on the increase of pressure at constant volume between 0° and 100° Cent. It gives most satisfactory results for expansion at constant pressure, compression at constant temperature, and also (I think) for cooling by free expansion" [i. e. the cooling effect in our experiments].

"Carbonic Acid Gas.

P pressure in pounds per square foot.

V volume of one pound in cubic feet.

 P_{o} one atmosphere.

V₀ theoretical volume, in the state of perfect gas, of one lb. at the pressure P₀ and the temperature of melting ice.

 P_0V_0 for carbonic acid 17116 feet, $\log P_0V_0 = 4.2334023$.

 $(P_0 V_0 \text{ actually, at } 0^{\circ}, 17145.)$

 K_p dynam. spec. heat at constant pressure 300.7 feet; $\log K_p = 2.4781334$.

C absolute temperature of melting ice, 274° Cent.

^{*} Philosophical Transactions, June 1853.

"The absolute zeros of gaseous tension and of heat are supposed sensibly to coincide, i. e. z is supposed inappreciably small.

" Formulæ:

$$\frac{PV}{P_0V_0} = \frac{T+C}{C} - \frac{a}{T+C} \frac{V_0}{V}. \qquad (1)$$

$$a = 1.9, \quad \log a = 0.2787536.$$

"Cooling by free expansion, supposing the perfect gas thermometer to give the true scale of absolute temperatures:

$$\delta T = \frac{P_0 V_0}{K_P} \cdot \frac{3a}{T + C} \left\{ \frac{V_0}{V_1} - \frac{V_0}{V_2} \right\} . \qquad (2)^*$$

$$\log \frac{3P_0 V_0 a}{K_P} = 2.5111438.$$

By substituting for $\frac{V_0}{V_1}$ and $\frac{V_0}{V_2}$ their approximate values $\frac{C}{T+C} \cdot \frac{P_1}{P_0}$ and $\frac{C}{T+C} \cdot \frac{P_2}{P_0}$, we reduce it to

$$\delta = \frac{3P_0V_0aC}{K_P(T+C)^2},$$

from which we have calculated the theoretical results for different temperatures shown above, which agree remarkably well with those we have obtained from observation.

The interpretation given above for the experimental results on mixtures of carbonic acid and air depends on the assumption (rendered probable as a very close approximation to the truth, by Dalton's law), that in a mixture each gas retains all its physical properties unchanged by the presence of the other. This assumption, however, may be only approximately true, perhaps similar in accuracy to Boyle's and Gaylussac's laws of compression and expansion by heat; and the theory of gases would be very much advanced by accurate comparative experiments on all the physical properties of mixtures and of their components separately. Towards this object we have experimented on the thermal effect of the mutual interpenetration of carbonic acid and air. In one experiment we found that when 7500 cubic inches of carbonic acid at the atmospheric pressure were mixed with 1000 cubic inches of common air and a perfect mutual interpenetration had taken place, the temperature had fallen by about '2° Cent. We intend to try more exact experiments on this subject.

THEORETICAL DEDUCTIONS.

Section I. On the Relation between the Heat evolved and the Work spent in Compressing a Gas kept at constant temperature.

This relation is not a relation of simple mechanical equivalence, as was supposed by MAYER in his 'Bemerkungen ueber die Kräfte der Unbelebten Natur,'

* Obtained by using Mr. Rankine's formula (1) in the general expression for the cooling effect given in our former paper, and repeated below as equation (15) of Section V.

2 x

† Annalen of Wöhler and Liebig, May 1842.

in which he founded on it an attempt to evaluate numerically the mechanical equivalent of the thermal unit. The heat evolved may be less than, equal to, or greater than the equivalent of the work spent, according as the work produces other effects in the fluid than heat, produces only heat, or is assisted by molecular forces in generating heat, and according to the quantity of heat, greater than, equal to, or less than that held by the fluid in its primitive condition, which it must hold to keep itself at the same temperature when compressed. The $\dot{\alpha}$ priori assumption of equivalence, for the case of air, without some special reason from theory or experiment, is not less unwarrantable than for the case of any fluid whatever subjected to compression. Yet it may be demonstrated* that water below its temperature of maximum density (39° 1 FAHR.), instead of evolving any heat at all when compressed, actually absorbs heat, and at higher temperatures evolves heat in greater or less, but probably always very small, proportion to the equivalent of the work spent; while air, as will be shown presently, evolves always, at least when kept at any temperature between 0° and 100° Cent., somewhat more heat than the work spent in compressing it could alone create. The first attempts to determine the relation in question, for the case of air, established an approximate equivalence without deciding how close it might be, or the direction of the discrepance, if any. Thus experiments "On the Changes of Temperature produced by the Rarefaction and Condensation of Air **," showed an approximate agreement between the heat evolved by compressing air into a strong copper vessel under water, and the heat generated by an equal expenditure of work in stirring a liquid; and again, conversely, an approximate compensation of the cold of expansion when air in expanding spends all its work in stirring its own mass by rushing through the narrow passage of a slightly opened stopcock. Again, theory;, without any doubtful hypothesis, showed from REGNAULT's observations on the pressure and latent heat of steam, that unless the density of saturated steam differs very much from what it would be if following the gaseous laws of expansion and compression, the heat evolved by the compression of air must be sensibly less than the equivalent of the work spent when the temperature is as low as 0° Cent., and very considerably greater than that equivalent when the temperature is above 40° or 50°. Mr. RANKINE is, so far as we know, the only other writer who independently admitted the necessity of experiment on the subject, and he was probably not aware of the experiments which had been made in 1844, on the rarefaction and condensation of air, when he remarked \(\), that "the value of z is

^{*} Dynamical Theory of Heat, § 63, equation (b.), Trans. Roy. Soc. Edinb. vol. xvi. p. 290; or Phil. Mag. vol. iv. Series 4. p. 425.

[†] Communicated to the Royal Society, June 20, 1844, and published in the Philosophical Magazine, May 1845.

[‡] Appendix to "Account of Carnor's Theory," Roy. Soc. Edinburgh, April 30, 1849, Transactions, vol. xvi. p. 568; confirmed in the Dynamical Theory, § 22, Transactions Roy. Soc. Edinb. March 17, 1851; and Phil. Mag. vol. iv. Series 4. p. 20.

[§] Mechanical Action of Heat, Section II. (10.), communicated to the Roy. Soc Edinb. Feb. 4, 1850, ransactions, vol. xx. p. 166.

unknown; and as yet no experimental data exist by which it can be determined" (z denoting in his expressions a quantity the vanishing of which for any gas would involve the equivalence in question). In further observing that probably z is small in comparison with the reciprocal of the coefficient of expansion, Mr. Rankine virtually adopted the equivalence as probably approximate; but in his article "On the Thermic Phenomena of Currents of Elastic Fluids*," he took the first opportunity of testing it closely, afforded by our preliminary experiments on the thermal effects of air escaping through narrow passages.

We are now able to give much more precise answers to the question regarding the heat of compression, and to others which rise from it, than those preliminary experiments enabled us to do. Thus if K denote the specific heat under constant pressure, of air or any other gas, issuing from the plug in the experiments described above, the quantity of heat that would have to be supplied, per pound of the fluid passing, to make the issuing stream have the temperature of the bath, would be $K\delta$, or

$$Km\frac{(P-P')}{\Pi}$$
,

where m is equal to '26° for air and 1'15° for carbonic acid, since we found that the cooling effect was simply proportional to the difference of pressure in each case, and was '0176° per pound per square inch, or '26 per atmosphere, for air, and about $4\frac{1}{2}$ times as much for carbonic acid. This shows precisely how much the heat of friction in the plug falls short of compensating the cold of expansion. But the heat of friction is the thermal equivalent of all the work done actually in the narrow passages by the air expanding as it flows through. Now this, in the cases of air and carbonic acid, is really not as much as the whole work of expansion, on account of the deviation from Boyle's law to which these gases are subject; but it exceeds the whole work of expansion in the case of hydrogen which presents a contrary deviation; since P'V', the work which a pound of air must do to escape against the atmospheric pressure, is, for the two former gases, rather greater, and for hydrogen rather less, than PV. which is the work done on it in pushing it through the spiral up to the plug. In any case, w denoting the whole work of expansion, w-(P'V'-PV) will be the work actually spent in friction within the plug; and

$$\frac{1}{J} \{ w - (P'V' - PV) \}$$

will be the quantity of heat into which it is converted, a quantity which, in the cases of air and carbonic acid, falls short by

$$Km\frac{P-P'}{\Pi}$$

of compensating the cold of expansion. If therefore H denote the quantity of heat

^{*} Mechanical Action of Heat, Subsection 4, communicated to the Roy. Soc. Edinb. Jan. 4, 1853, Transactions, vol. xx. p. 580.

that would exactly compensate the cold of expansion, or which amounts to the same, the quantity of heat that would be evolved by compressing a pound of the gas from the volume V' to the volume V, when kept at a constant temperature, we have

$$\frac{1}{J}\{w - (P'V' - PV)\} = H - Km\frac{P - P'}{\Pi},$$

whence

$$\mathbf{H} = \frac{w}{\mathbf{J}} + \left\{ -\frac{1}{\mathbf{J}} (\mathbf{P}'\mathbf{V}' - \mathbf{P}\mathbf{V}) + \mathbf{K}m \frac{\mathbf{P} - \mathbf{P}'}{\mathbf{\Pi}} \right\}$$

Now, from the results derived by REGNAULT from his experiments on the compressibility of air, of carbonic acid, and of hydrogen, at three or four degrees above the freezing-point, we find, approximately,

$$\frac{P'V'-PV}{PV}=f\frac{P-P'}{\Pi},$$

where

f = .00082 for air,

f= .0064 for carbonic acid,

and

f = -.00043 for hydrogen.

No doubt the deviations from Boyle's law will be somewhat different at the higher temperature (about 15° or 16° Cent.) of the bath in our experiments, probably a little smaller for air and carbonic acid, and possibly greater for hydrogen; but the preceding formula may express them accurately enough for the rough estimate which we are now attempting.

We have, therefore, for air or carbonic acid,

$$\mathbf{H} = \frac{w}{\mathbf{J}} + \left(\mathbf{K}m - \frac{\mathbf{P}Vf}{\mathbf{J}}\right) \frac{\mathbf{P} - \mathbf{P}'}{\mathbf{\Pi}} = \frac{w}{\mathbf{J}} + \frac{\mathbf{P}V}{\mathbf{J}} \left(\frac{\mathbf{J}Km}{\mathbf{P}V} - f\right) \frac{\mathbf{P} - \mathbf{P}'}{\mathbf{\Pi}}.$$

The values of JK and PV for the three gases in the circumstances of the experiments are as follow:—

For atmospheric air JK= $1390 \times 238 = 331$ For carbonic acid JK= $1390 \times 217 = 301$ For hydrogen . . JK= $1390 \times 3.4046 = 4732$

and for atmospheric air, at 15° Cent. PV=26224(1+15×00366)=27663

for carbonic acid, at 10° Cent. $PV = 17154(1+10 \times 00366) = 17782$

for hydrogen . . at 10° Cent. PV= $378960(1+10\times 00367)=393000$.

Hence we have, for air and carbonic acid,

$$\mathbf{H} = \frac{w}{\mathbf{J}} + \frac{\mathbf{PV}}{\mathbf{J}} \cdot \lambda \frac{\mathbf{P} - \mathbf{P}'}{\mathbf{\Pi}},$$

where λ denotes '0024 for air, and '013 for carbonic acid; showing (since these values of λ are positive) that in the case of each of these gases, more heat is evolved in compressing it than the equivalent of the work spent (a conclusion that would hold for hydrogen even if no cooling effect, or a heating effect less than a certain limit, were observed for it in our form of experiment). To find the proportion which this excess bears to the whole heat evolved, or to the thermal equivalent of the work spent

in the compression, we may use the expression

$$w = PV \log \frac{P}{P'}$$

as approximately equal to the mechanical value of either of those energies; and we thus find for the proportionate excess,

$$\frac{H - \frac{1}{J}w}{\frac{1}{J}w} = \lambda \frac{P - P'}{\Pi \log \frac{P}{P'}} = 0024 \frac{P - P'}{\Pi \log \frac{P}{P'}} \text{ for air,}$$

$$= 013 \frac{P - P'}{\Pi \log \frac{P}{P'}} \text{ for carbonic acid.}$$

or

This equation shows in what proportion the heat evolved exceeds the equivalent of the work spent in any particular case of compression of either gas. Thus for a very small compression from $P'=\Pi$, the atmospheric pressure, we have

$$\log \frac{P}{P'} = \log \left(1 + \frac{P - \Pi}{\Pi} \right) = \frac{P - \Pi}{\Pi} \text{ approximately,}$$

$$\frac{H - \frac{1}{J}w}{\frac{1}{J}w} = 0024 \text{ for air,}$$

and therefore

or

= 013 for carbonic acid.

Therefore, when slightly compressed from the ordinary atmospheric pressure, and kept at a temperature of about 60° Fahr., common air evolves more heat by $\frac{1}{417}$, and carbonic acid more by $\frac{1}{77}$ than the amount mechanically equivalent to the work of compression. For considerable compressions from the atmospheric pressure, the proportionate excesses of the heat evolved are greater than these values, in the ratio of the Napierian logarithm of the number of times the pressure is increased, to this number diminished by 1. Thus, if either gas be compressed from the standard state to double density, the heat evolved exceeds the thermal equivalent of the work spent, by $\frac{1}{290}$ in the case of air, and by $\frac{1}{53}$ in the case of carbonic acid.

As regards these two gases, it appears that the observed cooling effect was chiefly due to an actual preponderance of the mechanical equivalent of the heat required to compensate the cold of expansion over the work of expansion, but that rather more than one-fourth of it in the case of air, and about one-third of it in the case of carbonic acid, depended on a portion of the work of expansion going to do the extra work spent by the gas in issuing against the atmospheric pressure above that gained by it in being sent into the plug. On the other hand, in the case of hydrogen, in such an experiment as we have performed, there would be a heating effect, if the

work of expansion were precisely equal to the mechanical equivalent of the cold of expansion, since not only the whole work of expansion, but also the excess of the work done in forcing the gas in above that performed by it in escaping, is spent in friction in the plug. Since we have observed actually a cooling effect, it follows that the heat absorbed in expansion must exceed the equivalent of the work of expansion, enough to over-compensate the whole heat of friction mechanically equivalent, as this is, to the work of expansion together with the extra work of sending the gas into the plug above that which it does in escaping. In the actual experiment* we found a cooling effect of '076°, with a difference of pressures, P—P', equal to 53'7 lbs. per square inch, or 3'7 atmospheres. Now the mechanical value of the specific heat of a pound of hydrogen is, according to the result stated above, 4732 foot-pounds, and hence the mechanical value of the heat that would compensate the observed cooling effect per pound of hydrogen passing is 360 foot-pounds. But, according to Regnault's experiments on the compression of hydrogen, quoted above, we have

$$PV-P'V'=PV\times 00043\frac{P-P'}{\Pi}$$
 approximately;

and as the temperature was about 10° in our experiment, we have, as stated above, PV=393000.

Hence, for the case of the experiment in which the difference of pressures was 3.7 atmospheres, or

 $\frac{P-P'}{\Pi}=3.7,$

we have

$$PV - P'V' = 625$$
:

that is, 625 foot-pounds more of work, per pound of hydrogen, is spent in sending the hydrogen into the plug at 4.7 atmospheres of pressure, than would be gained in allowing it to escape at the same temperature against the atmospheric pressure. Hence the heat required to compensate the cold of expansion, is generated by friction from (1) the actual work of expansion, together with (2) the extra work of 625 foot-pounds per pound of gas, and (3) the amount equivalent to 360 foot-pounds which would have to be communicated from without to do away with the residual cooling effect observed. Its mechanical equivalent therefore exceeds the work of expansion by 985 foot-pounds; which is $\frac{1}{630}$ of its own amount, since the work of expansion in the circumstances is approximately $393000 \times \log 4.7 = 608000$ foot-pounds. Conversely, the heat evolved by the compression of hydrogen at 10° Cent., from 1 to 4.7 atmospheres, exceeds by $\frac{1}{630}$ the work spent. The corresponding excess in the case

^{*} From the single experiment we have made on hydrogen we cannot conclude that at other pressures a cooling effect proportional to the difference of pressures would be observed, and therefore we confine the comparison of the three gases to the particular pressure used in the hydrogen experiment. It should be remarked too, that we feel little confidence in the value assigned to the thermal effect for the case observed in the experiment on hydrogen, and only consider it established that it is a cooling effect, and very small.

of atmospheric air, according to the result obtained above, is $\frac{1}{174}$, and in the case of carbonic acid $\frac{1}{32}$.

It is important to observe how much less close is the compensation in carbonic acid than in either of the other gases, and it appears probable that the more a gas deviates from the gaseous laws, or the more it approaches the condition of a vapour at saturation, the wider will be the discrepancy. We hope, with a view to investigating further the physical properties of gases, to extend our method of experimenting to steam (which will probably present a large cooling effect), and perhaps to some other vapours.

In Mr. Joule's original experiment* to test the relation between heat evolved and work spent in the compression of air, without an independent determination of the mechanical equivalent of the thermal unit, air was allowed to expand through the aperture of an open stopcock from one copper vessel into another previously exhausted by an air-pump, and the whole external thermal effect on the metal of the vessels, and a mass of water below which they are kept, was examined. We may now estimate the actual amount of that external thermal effect, which observation only showed to be insensibly small. In the first place it is to be remarked, that, however the equilibrium of pressure and temperature is established between the two air vessels, provided only no appreciable amount of work is emitted in sound, the same quantity of heat must be absorbed from the copper and water to reduce them to their primitive temperature; and that this quantity, as was shown above, is equal to

$$\frac{\text{PV}}{\text{J}} \times .0024 \times \frac{\text{P} - \text{P'}}{\Pi} = \frac{27000 \times .0024}{1390} \times \frac{\text{P} - \text{P'}}{\Pi} = .046 \frac{\text{P} - \text{P'}}{\Pi}.$$

In the actual experiments the exhausted vessel was equal in capacity to the charged vessel, and the latter contained '13 of a pound of air under 21 atmospheres of pressure, at the commencement. Hence $P' = \frac{1}{2} P$, and

$$\frac{P-P'}{\Pi} = 10.5$$
;

and the quantity of heat required from without to compensate the total internal cooling effect must have been

$$.046 \times 10.5 \times .13 = .063$$
.

This amount of heat, taken from $16\frac{1}{2}$ lbs. of water, 28 lbs. of copper, and 7 lbs. of tinned iron, as in the actual experiment, would produce a lowering of temperature of only 003° Cent. We need not therefore wonder that no sensible external thermal effect was the result of the experiment when the two copper vessels and the pipe connecting them were kept under water, stirred about through the whole space

^{*} The second experiment mentioned in the abstract published in the Proceedings of the Royal Society, June 20, 1844, and described in the Philosophical Magazine, May 1845, p. 377.

surrounding them, and that similar experiments, more recently made by M. Regnault, should have led only to the same negative conclusion.

If, on the other hand, the air were neither allowed to take in heat from nor to part with heat to the surrounding matter in any part of the apparatus, it would experience a resultant cooling effect (after arriving at a state of uniformity of temperature as well as pressure) to be calculated by dividing the preceding expression for the quantity of heat which would be required to compensate it, by 17, the specific heat of air under constant pressure. The cooling effect on the air itself therefore amounts to

$$0^{\circ}.27 \times \frac{P-P'}{\Pi},*$$

which is equal to 2°·8, for air expanding, as in Mr. Joule's experiment, from 21 atmospheres to half that pressure, and is 900 times as great as the thermometric effect when spread over the water and copper of the apparatus. Hence our present system, in which the thermometric effect on the air itself is directly observed, affords a test hundreds of times more sensitive than the method first adopted by Mr. Joule, and no doubt also than that recently practised by M. Regnault, in which the dimensions of the various parts of the apparatus (although not yet published) must have been on a corresponding scale, or in somewhat similar proportions, to those used formerly by Mr. Joule.

Section II. On the Density of Saturated Steam.

The relation between the heat evolved and the work spent, approximately established by the air-experiments communicated to the Royal Society in 1844, was subjected to an independent indirect test by an application of Carnot's theory, with values of "Carnot's function" which had been calculated from Regnault's data as to the pressure and latent heat of steam, and the assumption (in want of experimental data), that the density varies according to the gaseous laws. The verification thus obtained was very striking, showing an exact agreement with the relation of equivalence at a temperature a little above that of observation, and an agreement with the actual experimental results quite within the limits of the errors of observation; but a very wide discrepancy from equivalence for other temperatures. The following Table is extracted from the Appendix to the "Account of Carnot's Theory" in which the theoretical comparison was first made, to facilitate a comparison with what we now know to be the true circumstances of the case.

* It is worthy of remark that this, the expression for the cooling effect experienced by a mass of atmospheric air expanding from a bulk in which its pressure is P to a bulk in which, at the same (or very nearly the same) temperature its pressure is P', and spending all its work of expansion in friction among its own particles, agrees very closely with the expression, $26 \times \frac{P-P'}{\Pi}$, for the cooling effect in the somewhat different circumstances of our experiments.