

**Exact solution of Schrödinger equation
in the case of reduction to *Riccati* type of ODE.**

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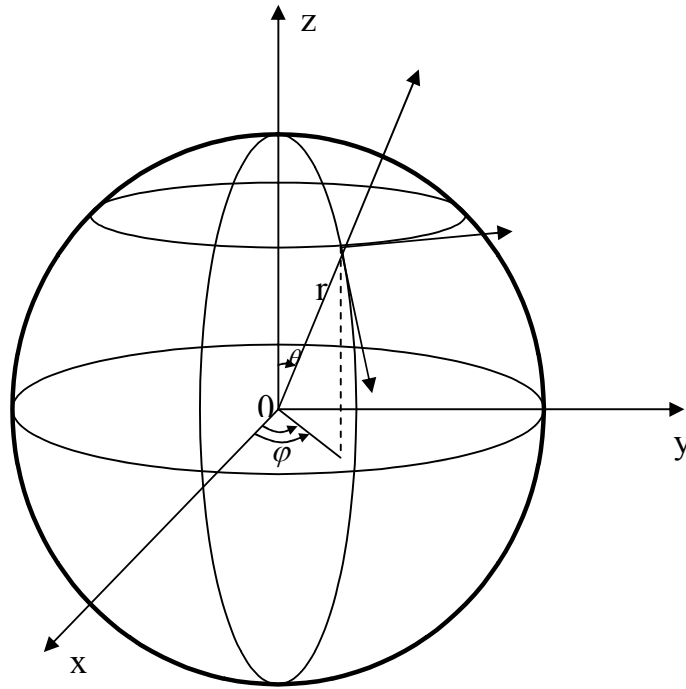
Here is presented a new type of exact solution of Schrödinger equation in the case of its reduction to *Riccati* type of ordinary differential equations.

Due to a very special character of *Riccati*'s type equation, its general solution is proved to have a *proper gap* of components of the particle wavefunction (which is known to be determining a proper quantum state of the particle).

It means a possibility of sudden transformation *or transmutation* of quantum state of the particle (*from one meaning of wavefunction to another*), at definite moment of parametrical time.

Besides, in the case of *spherical symmetry* of particle potential V in position space, as well as *spherical symmetry* of quantum system E total energy, such a solution is proved to be a multiplying of *Bessel* function (for radial component) & *Legendre* spherical function (for angle component), in spherical coordinate system.

In accordance with [1-4], Schrödinger equation describes the quantum state of a physical system changing in time or how the wavefunction of a proper particle evolves over time. We observe below only the time-independent solutions ($\partial/\partial t = 0$) of Schrödinger equation [5-6], for the reason that in case of time-independent particle potential V in position space, the general solutions of Schrödinger equation could be obtained from time-independent solutions by multiplying it on a proper time-coefficient $e^{-\frac{i}{\hbar} \cdot E \cdot t}$, where E – is a total energy of quantum system, $i = \sqrt{-1}$ – is the imaginary unit, \hbar – is h -bar.



Pic.1. Spherical coordinate system R, θ, φ .

Schrödinger equation for a single particle in the 3-dimensional case in the presence of a potential V , should be represented [5] in a spherical coordinate system R, θ, φ (Pic.1):

$$\frac{\hbar^2}{2m} \Delta \psi = U \psi, \quad (1.1)$$

- here m - is the mass of particle, $m = const \neq 0$; ψ - the wavefunction of particle in position space, which is known to be the most complete description of particle quantum mechanical behavior, $\psi = \psi(R, \theta, \varphi)$; $U = (V - E)$, where V - is the potential of the particle in position space, E - is a quantum system total energy, $U = U(R, \theta, \varphi)$; \hbar - is the reduced Planck constant.

Besides, in spherical coordinate system:

$$\Delta\psi = \frac{\partial^2\psi}{\partial R^2} + \frac{2}{R} \frac{\partial\psi}{\partial R} + \frac{1}{R^2 \sin^2\theta} \frac{\partial^2\psi}{\partial \varphi^2} + \frac{1}{R^2} \frac{\partial^2\psi}{\partial \theta^2} + \frac{1}{R^2} \text{ctg}\theta \frac{\partial\psi}{\partial \theta} .$$

Also it should be taken into consideration that the *square of wavefunction* ψ is known to be a *proper probability amplitude* to detect such a particle in position space (*per unit of volume of such a space*).

In accordance with [5-6], we should find an exact solution of above equation (1.1) in a form below:

$$\begin{aligned} \psi(R, \theta, \varphi) &= \psi(R) \cdot \psi(\theta) \cdot \psi(\varphi), \\ U(R, \theta, \varphi) &= U(R) \cdot U(\theta) \cdot U(\varphi) . \end{aligned} \quad (1.2)$$

Having substituted (1.2) into (1.1), we obtain:

$$\begin{aligned} \frac{1}{\psi(R)} \cdot \frac{d^2\psi(R)}{dR^2} + \frac{2}{R} \cdot \left(\frac{1}{\psi(R)} \right) \cdot \frac{d\psi(R)}{dR} + \frac{1}{R^2 \sin^2\theta} \cdot \left(\frac{1}{\psi(\varphi)} \right) \cdot \frac{d^2\psi(\varphi)}{d\varphi^2} + \\ + \frac{1}{R^2} \cdot \left(\frac{1}{\psi(\theta)} \right) \cdot \frac{d^2\psi(\theta)}{d\theta^2} + \frac{\text{ctg}\theta}{R^2} \cdot \left(\frac{1}{\psi(\theta)} \right) \cdot \frac{d\psi(\theta)}{d\theta} = \frac{2m}{\hbar^2} \cdot U(R) \cdot U(\theta) \cdot U(\varphi) . \end{aligned}$$

Let's make a proper replacement of variables $v = \psi' / \psi$ over all the arguments R, θ, φ :

$$\begin{aligned} \left(v^2(R) + v'(R) \right) + \frac{2}{R} \cdot v(R) + \frac{1}{R^2 \cdot \sin^2\theta} \cdot \left(v^2(\varphi) + v'(\varphi) \right) + \\ + \frac{1}{R^2} \cdot \left(v^2(\theta) + v'(\theta) \right) + \frac{\text{ctg}\theta}{R^2} \cdot v(\theta) = \frac{2m}{\hbar^2} \cdot U(R) \cdot U(\theta) \cdot U(\varphi) \end{aligned} \quad (1.3)$$

The main idea is to find the solution of (1.1) like *Riccati's type* [7], that's why let's make an additional assumption below:

$$\begin{aligned} (v^2(R) + v'(R)) &= f_1(R), \\ (v^2(\theta) + v'(\theta)) &= f_2(\theta), \\ (v^2(\varphi) + v'(\varphi)) &= f_3(\varphi). \end{aligned} \tag{1.4}$$

Having substituted (1.4) into (1.3), we obtain an equation below:

$$\begin{aligned} f_1(R) + \frac{2}{R} \cdot v(R) + \frac{I}{R^2 \cdot \sin^2 \theta} \cdot f_3(\varphi) + \\ + \frac{I}{R^2} \cdot f_2(\theta) + \frac{\text{ctg} \theta}{R^2} \cdot v(\theta) = \frac{2m}{\hbar^2} \cdot U(R) \cdot U(\theta) \cdot U(\varphi). \end{aligned}$$

According to [6], we will present a function $U = (V - E)$ in the axe-symmetrical form. It means that $\partial U / \partial \varphi = 0$

$$U(\varphi) = \text{const} = C_0 \neq 0.$$

Thus, from the last equation we obtain (*in the axe-symmetrical case*):

$$\begin{aligned} f_1(R) + \frac{2}{R} \cdot v(R) + \frac{I}{R^2 \cdot \sin^2 \theta} \cdot f_3(\varphi) + \\ + \frac{I}{R^2} \cdot f_2(\theta) + \frac{\text{ctg} \theta}{R^2} \cdot v(\theta) = \frac{2m \cdot C_0}{\hbar^2} \cdot U(R) \cdot U(\theta). \end{aligned} \tag{1.5}$$

To separate a proper variables in (1.5), we should make an assumption (*below $C_1 \neq 0$*):

$$f_3(\varphi) = \text{const} = C_1 \tag{1.6}$$

Above (1.6) leads us to equality below:

$$f_1(R) + \frac{2}{R} \cdot v(R) + \frac{C_1}{R^2 \cdot \sin^2 \theta} + \frac{I}{R^2} \cdot f_2(\theta) + \frac{\text{ctg} \theta}{R^2} \cdot v(\theta) = \frac{2m \cdot C_0}{\hbar^2} \cdot U(R) \cdot U(\theta),$$

- or:

$$U(\theta) \cdot \left(\frac{2m \cdot C_0}{\hbar^2} \right) \cdot U(R) \cdot R^2 - \left(R^2 \cdot f_1(R) + 2R \cdot v(R) \right) = \frac{C_1}{\sin^2 \theta} + f_2(\theta) + \text{ctg} \theta \cdot v(\theta) .$$

From equality above we conclude that variables in (1.5) could be separated only in cases:

- 1) $U(\theta) = \text{const} = C_2 \neq 0$;
- 2) $U(R) \cdot R^2 = \text{const} = C_3 \neq 0$.

Above 2) case was considered in [5]. Besides, as for the expression $U(R) = C_3 / R^2$ itself: presense of it means that meaning of the particle potential V is proved to differ from a total energy of quantum system E (at given locus of position space) on a proper term $\sim (1/R^2)$, which could be associated with a proper influence on a particle the field of centrifugal forces.

Let's consider above 1) case: as for the condition $U(\theta) = \text{const} = C_2$, it means we explore the case of stationary, *spherical symmetry* of particle potential V , as well as case of stationary, *spherical symmetry* of total energy of quantum system E .

In such an assumptions, last equation could be represented in form below:

$$\frac{2m \cdot C_0 \cdot C_2}{\hbar^2} \cdot U(R) \cdot R^2 - \left(R^2 \cdot f_1(R) + 2R \cdot v(R) \right) = \frac{C_1}{\sin^2 \theta} + f_2(\theta) + \text{ctg} \theta \cdot v(\theta) ,$$

- or

$$-R^2 \cdot \left(-\frac{2m \cdot C_0 \cdot C_2}{\hbar^2} \cdot U(R) + f_1(R) \right) - 2R \cdot v(R) = \frac{C_1}{\sin^2 \theta} + f_2(\theta) + \text{ctg} \theta \cdot v(\theta) .$$

To separate a proper variables (R , θ) in above equality, we should make a conclusion

$$-R^2 \cdot \left(-\frac{2m \cdot C_0 \cdot C_2}{\hbar^2} \cdot U(R) + f_1(R) \right) - 2R \cdot v(R) = \text{const} = \tau \tag{1.7}$$

$$\frac{C_1}{\sin^2 \theta} + f_2(\theta) + \text{ctg} \theta \cdot v(\theta) = \text{const} = \tau .$$

If we determine the functions $f_1(R)$ & $f_2(\theta)$ from an equalities (1.7), then substituting of such a functions into (1.4), we should obtain:

$$\begin{aligned} \left(v^2(R) + v'(R)\right) &= f_1(R) = -\frac{\tau}{R^2} - \frac{2v(R)}{R} + \frac{2m \cdot C_0 \cdot C_2}{\hbar^2} \cdot U(R), \\ \left(v^2(\theta) + v'(\theta)\right) &= f_2(\theta) = \tau - \text{ctg } \theta \cdot v(\theta) - \frac{C_1}{\sin^2 \theta}, \end{aligned}$$

- or

$$\begin{aligned} v'(R) &= -v^2(R) - \frac{2}{R} \cdot v(R) - \left(\frac{\tau}{R^2} - \frac{2m \cdot C_0 \cdot C_2}{\hbar^2} \cdot U(R) \right), \\ v'(\theta) &= -v^2(\theta) - \text{ctg } \theta \cdot v(\theta) + \left(\tau - \frac{C_1}{\sin^2 \theta} \right). \end{aligned} \tag{1.8}$$

In accordance with previous assumption, we obtain a proper system of *Riccati* equations [7] in regard to the functions $v = \psi' / \psi$ (which are functions of R , θ , respectively).

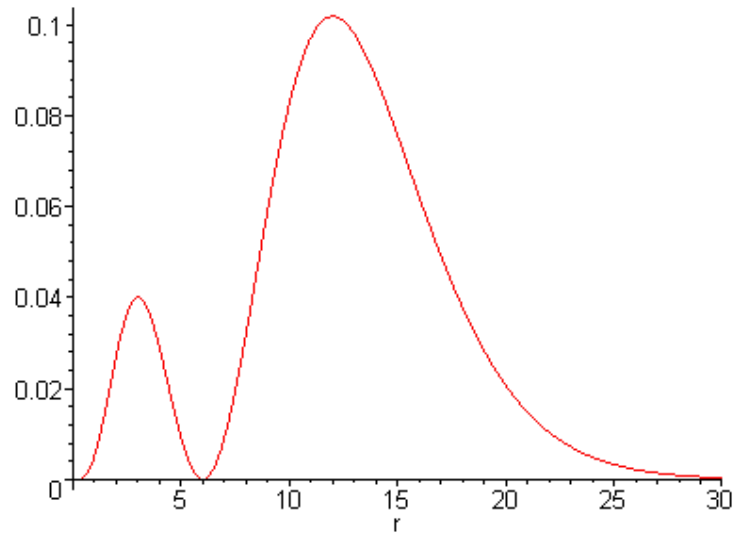
Due to a very special character of such an equations, it's general solution is known to have a *proper gap* of above components of such a functions $v = \psi' / \psi$ [7].

It means a possibility of sudden transformation *or transmutation* of quantum mechanical state of the particle (from one meaning of wavefunction to another) at definite moment of parametrical time, or existence of such a continuous solution only at some definite, restricted range of parameter t [8].

Expressing of functions ψ from equality $v = \psi' / \psi$ in regard to R , θ in (1.8), we obtain:

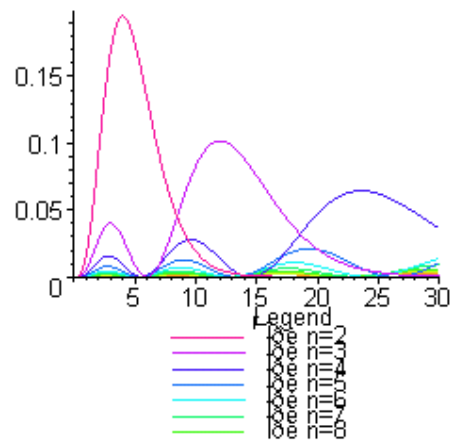
$$\begin{aligned} \psi''(R) + \frac{2}{R} \cdot \psi'(R) + \left(\frac{\tau}{R^2} - \frac{2m \cdot C_0 \cdot C_2}{\hbar^2} \cdot U(R) \right) \cdot \psi(R) &= 0, \\ \psi''(\theta) + \text{ctg } \theta \cdot \psi'(\theta) + \left(\frac{C_1}{\sin^2 \theta} - \tau \right) \cdot \psi(\theta) &= 0. \end{aligned} \tag{1.9}$$

The 1-st of equations (1.9) ($\tau \neq 0, U \neq 0$) – is known to be *the Bessel* ordinary differential equation, it's solutions are a proper *Bessel* functions (Pic.2; besides, see for example 2.162, 1a [7]):



Pic.2. Bessel functions.

The 2-nd equation of (1.9) – is known to be *the Legendre* ordinary differential equation, it's solutions are a proper *Legendre* spherical functions (Pic.3; see for example 2.240 [7]):



Pic.3. Legendre (spherical) functions.

Besides, we could make a conclusion that for function $U(R) = V(R) - E(R)$ are valid an equalities below:

$$V(R) \sim R^{m-2}, \quad E(R) \sim R^{-2}, \quad (1.10)$$

- where $m \neq 0$. Such an assumption has a real physical sense: indeed, potential V very often has a form $V(R) \sim 1/R$ (as in the case of gravitational field or case of electric charge field distribution), but a total energy E corresponds to the case of a proper set of discrete energy levels, which are known to be reproduced by harmonic oscillator: $E(R) \sim 1/R^2$.

Besides, in accordance to [7] solutions of 1-st of (1.9) could be expressed in terms of elementary transcendental functions only in case if $U(R) = 0$ or when $(n = 0, 1, 2, \dots)$:

$$\tau = 1/4 - (1/4) \cdot (2n+1)^2$$

Let's choose as a basic model the case of accordance of total energy of quantum system E to a proper set of discrete energy levels, which are known to be reproduced by a harmonic oscillator (in spherical coordinate system):

$$E(R) = \frac{((2n+1)^2 - 1)}{4R^2},$$

- also let's choose a constants C_0 & C_2 as below

$$\left(\frac{2m \cdot C_0 \cdot C_2}{\hbar^2} \right) = 1,$$

- then we obtain from (1.9)

$$\psi''(R) + \left(\frac{2}{R} \right) \cdot \psi'(R) - V(R) \cdot \psi(R) = 0,$$

- or, having chosen of a proper constant in representation of function $V(R)$, we obtain:

$$R^2 \cdot \psi''(R) + 2R \cdot \psi'(R) + R^m \cdot \psi(R) = 0,$$

- if we take into consideration (1.10).

The last equation is a classical case of *Bessel* ordinary differential equation [7] ($m \neq 0$):

$$\psi(R) = \left(\frac{1}{R^{1/2}} \right) \cdot Z_\nu \left(\frac{2}{m} \cdot R^{\frac{m}{2}} \right), \quad \nu = \frac{1}{m}.$$

If $m = 1$, then from above we obtain: $V(R) \sim -1/R$.

In such a case, solution of the 1-st equation of (1.9) should be represented as below:

$$\psi(R) = \left(\frac{1}{R^{1/2}} \right) \cdot Z_1(2\sqrt{R}).$$

As for the function $v = \psi'/\psi$ in regard to the argument φ , taking into consideration (1.4) & (1.6), we obtain [5-6]:

$$v'(\varphi) + v^2(\varphi) = C_1.$$

Besides, the requiring of limitation of such a solution for all $\varphi \in [0, 2\pi n]$, $n \in \mathbb{N}$, leads us to an equality below [5-6]:

$$\begin{aligned} \psi''(\varphi) - C_1 \cdot \psi(\varphi) &= 0, \\ \psi(\varphi) &= C_{01} \cdot \cos(\varphi \cdot \sqrt{|C_1|}) + C_{02} \cdot \sin(\varphi \cdot \sqrt{|C_1|}). \end{aligned}$$

- where $C_1 < 0$.

As mentioned above, we should take into consideration a requiring that the *square of wavefunction* ψ is to be a *proper probability amplitude* to detect such a particle in position space (*per unit of volume*).

It means that calculating of square of wavefunction ψ over all the volume of position space *must be equal to unit* = 1. Such a requiring is used to restrict the choosing of a constants of solution to be constructed. Besides, in the case of *divergence* of meanings of wavefunction ψ in calculating over the variables R , θ , φ , such a wavefunction ψ should be normalized to a proper *Dirac* δ -function [6].

References:

1. Schrödinger, Erwin (December 1926). "An Undulatory Theory of the Mechanics of Atoms and Molecules". *Phys. Rev.* 28 (6) **28**: 1049–1070.
[doi:10.1103/PhysRev.28.1049](https://doi.org/10.1103/PhysRev.28.1049).
2. Paul Adrien Maurice Dirac (1958). *The Principles of Quantum Mechanics* (4th ed.). Oxford University Press.
3. David J. Griffiths (2004). *Introduction to Quantum Mechanics* (2nd ed.). Benjamin Cummings. [ISBN 0131244051](#).
4. Serway, Moses, and Moyer (2004). *Modern Physics* (3rd ed.). Brooks Cole.
[ISBN 0534493408](#). See also:
<http://scienceworld.wolfram.com/physics/SchroedingerEquation.html>
5. Ershkov S.V. Self-similar solutions of Schroedinger equation. Part **I** // Moscow State University, proceedings of seminar on *Temporology exploring* (2005) (*in Russian*):
http://www.chronos.msu.ru/RREPORTS/yershkov_parametricheskaya.pdf.
6. Ershkov S.V. Self-similar solutions of Schroedinger equation. Part **II** // Moscow State University, proceedings of seminar on *Temporology exploring* (2005) (*in Russian*):
http://www.chronos.msu.ru/RREPORTS/yershkov_uravnenie.pdf.
7. Dr. E.Kamke. Hand-book for ordinary differential equations // Moscow: "Science" (1971).
8. Ershkov S.V., Schennikov V.V. Self-Similar Solutions to the Complete System of Navier-Stokes Equations for Axially Symmetric Swirling Viscous Compressible Gas Flow // *Comput. Math. and Math. Phys. J.* (2001) Vol.**41**, № 7. P.1117-1124.